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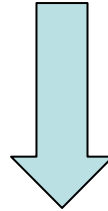
# **Pressure limit of hydrogen spontaneous ignition in a T-shaped channel**

Maxim Bragin, Dmitriy Makarov, Vladimir Molkov

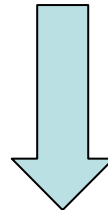
Hydrogen Safety Engineering and  
Research (HySAFER) Centre, University of Ulster

- ❖ Background
- ❖ Aim and objectives
- ❖ Experiment in T-shape channel (Golub et al., 2010)
- ❖ LES EDC model
- ❖ Modelling of membrane opening
  - ❖ Instant versus finite membrane opening
- ❖ Simulation results
  - ❖ Dynamics of spontaneous ignition at different storage pressure
- ❖ Conclusions

Pressure relief device (PRD) activation



Sudden release of hydrogen into air



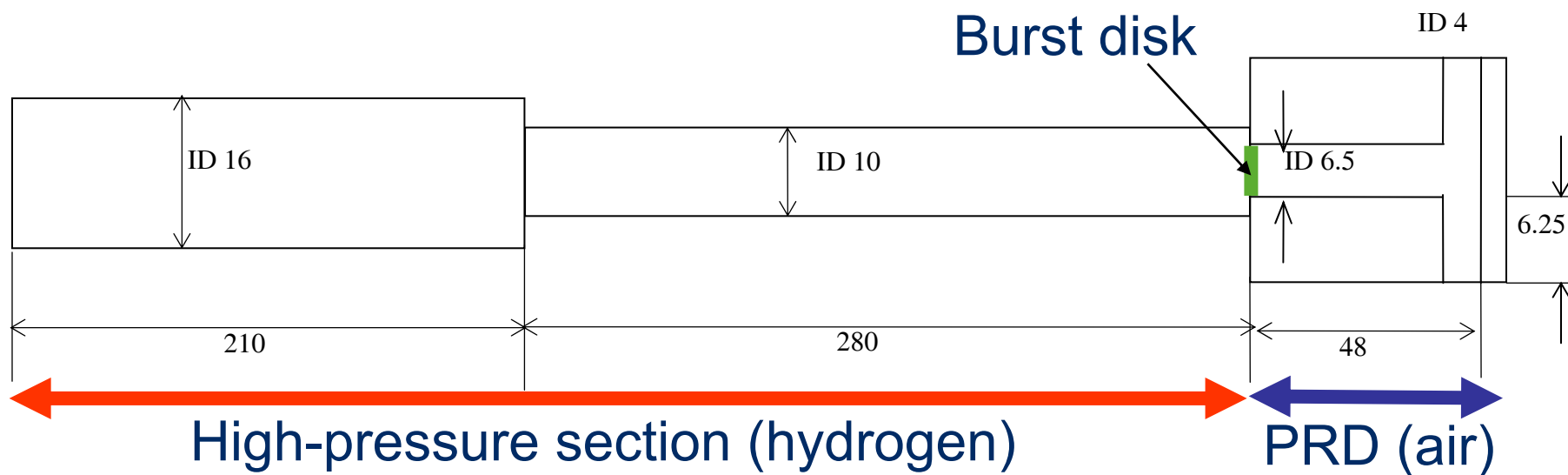
What is the lower pressure limit for spontaneous ignition in T-shape channel (PRD mock-up)?

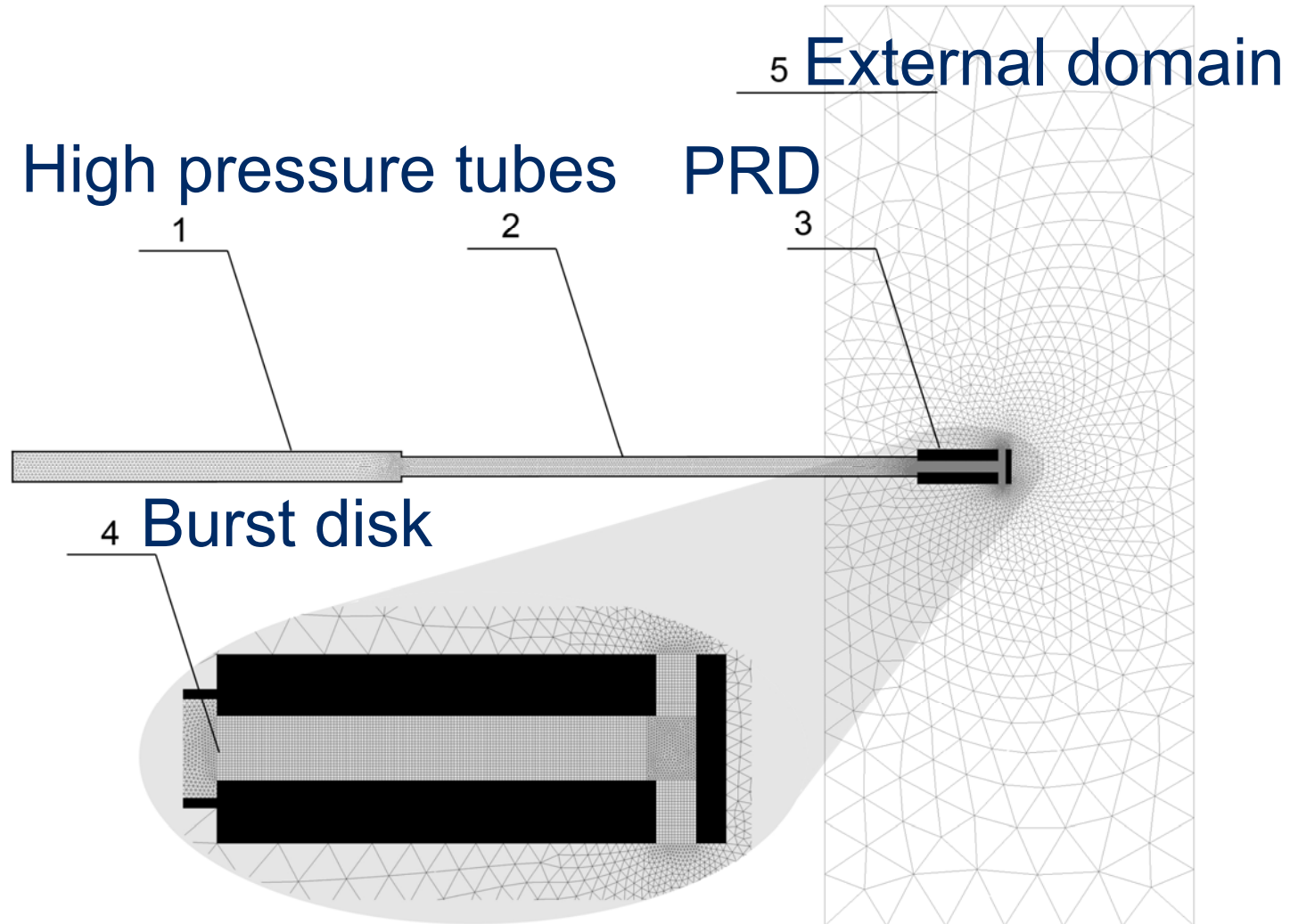
**Aim:** Develop contemporary model for predictive simulation of the phenomenon of hydrogen spontaneous ignition during sudden release from pressure relief devices (to be applied as a tool for hydrogen safety engineering)

**Objectives:**

- Develop Large Eddy Simulation (LES) model based on the Eddy Dissipation Concept (EDC), full chemistry, and modelling of a burst disk opening in a finite time
- Validate the LES EDC model and simulations against experimental data by Golub et al. (2010) on spontaneous ignition in T-shape channel (mock-up PRD)
- Understand dynamics of the ignition process and explain scattering in reported experimental data (Golub et al., 2010)
- Establish the lower pressure limit for spontaneous ignition in T-shape channel

Burst disk rupture at storage pressures from 13.5 to 29 bar. Reported: **no ignition at 13.5 bar**, and **spontaneous ignition at 29 bar** (private communication – ignition at **24.3 bar**)





Continuity equation: 
$$\frac{\partial \bar{\rho}}{\partial t} + \frac{\partial}{\partial x_j} (\bar{\rho} \tilde{u}_j) = 0$$

Momentum conservation equation:

$$\frac{\partial \bar{\rho} \tilde{u}_i}{\partial t} + \frac{\partial}{\partial x_j} (\bar{\rho} \tilde{u}_j \tilde{u}_i) = -\frac{\partial \bar{p}}{\partial x_i} + \frac{\partial}{\partial x_j} \left( \mu_{eff} \left( \frac{\partial \tilde{u}_i}{\partial x_j} + \frac{\partial \tilde{u}_j}{\partial x_i} - \frac{2}{3} \frac{\partial \tilde{u}_k}{\partial x_k} \delta_{ij} \right) \right) + \bar{\rho} g_i$$

Energy conservation equation:

$$\begin{aligned} \frac{\partial}{\partial t} (\bar{\rho} \tilde{E}) + \frac{\partial}{\partial x_j} (\tilde{u}_j (\bar{\rho} \tilde{E} + \bar{p})) = & \frac{\partial}{\partial x_j} \left[ \frac{\mu_{eff} c_p}{Pr_{eff}} \frac{\partial \tilde{T}}{\partial x_j} - \sum_m \tilde{h}_m \left( -\frac{\mu_{eff}}{Sc_{eff}} \frac{\partial \tilde{Y}_m}{\partial x_j} \right) \right] + \\ & + \tilde{u}_i \mu_{eff} \left( \frac{\partial \tilde{u}_i}{\partial x_j} + \frac{\partial \tilde{u}_j}{\partial x_i} - \frac{2}{3} \frac{\partial \tilde{u}_k}{\partial x_k} \delta_{ij} \right) \Big] + \sum_m R_m H_C \end{aligned}$$

Species conservation equation:

$$\frac{\partial}{\partial t} (\bar{\rho} \tilde{Y}_H) + \frac{\partial}{\partial x_j} (\bar{\rho} \tilde{u}_j \tilde{Y}_m) = \frac{\partial}{\partial x_j} \left( \frac{\mu_{eff}}{Sc_{eff}} \frac{\partial \tilde{Y}_m}{\partial x_j} \right) + R_m$$

- The effective viscosity in RNG SGS model is calculated as

$$\mu_{eff} = \mu \left[ 1 + H \left( \frac{\mu_s^2 \mu_{eff}}{\mu^3} - 100 \right) \right]^{1/3}, \text{ where}$$

$$\mu_s = \bar{\rho} \left( 0.157 V_{CV}^{1/3} \right)^2 \sqrt{2 \tilde{S}_{ij} \tilde{S}_{ij}}$$

$$S_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

- In highly turbulent flows,  $\mu_{eff} \approx \mu_s$  and the RNG model reduces to Smagorinsky model in laminar flow regions the argument of Heaviside function becomes negative and the model recovers molecular viscosity,  $\mu_{eff} = \mu$ .

- The effective Prandtl and Schmidt numbers are calculated following the RNG theory by Yakhot purely from the theoretical equation

$$\left| \frac{1/N_{eff} - 1.3929}{1/N - 1.3929} \right|^{0.6321} \left| \frac{1/N_{eff} + 2.3929}{1/N + 2.3929} \right|^{0.3679} = \frac{\mu}{\mu_{eff}}$$



- Eddy Dissipation Concept (EDC) gives an expression for a combustion rate based on an assumption that chemical reactions occur in **so-called fine structures** of Kolmogorov's scale, where the dissipation of turbulence energy takes place.
- In EDC model a source term in species transport equation is modelled as:

$$R_m = \frac{\rho(\xi^*)^2}{\tau^* [1 - (\xi^*)^3]} (Y_m^* - Y_m)$$

$R_m$  is a net rate of production of specie  $m$  by chemical reactions,

$\xi^*$  is the length fraction of the fine scale turbulent structures where the reaction occurs,

$Y_m^*$  is the fine scale species mass fraction (specie  $m$ ) after reacting over the time  $\tau^*$ ,

$Y_m$  is a species mass fraction for specie  $m$  in the surrounding fine scales state.

- The length fraction of fine structures:  $\xi^* = 2.1317u_\eta / u_{SGS}$   
 where  $u_{SGS} = \mu_t / (\rho \cdot L_{SGS})$  and  $L_{SGS} = 0.157V^{1/3}$   
 and the Kolmogorov's velocity  $u_\eta = \left( \frac{\mu \cdot u_{SGS}^3}{\rho \cdot L_{SGS}} \right)^{1/4}$
- Characteristic sub-grid eddy and Kolmogorov timescales are:  
 $\tau_{SGS} = L_{SGS} / u_{SGS}$  and  $\tau_\eta = \left( \frac{\mu \cdot L_{SGS}}{\rho \cdot u_{SGS}^3} \right)^{1/2}$
- The volume fraction of the fine scales is calculated as  $\xi^{*3}$   
 and species are assumed to react in the fine structures over a  
 time scale  $\tau^* = 0.4082\tau_\eta$ .

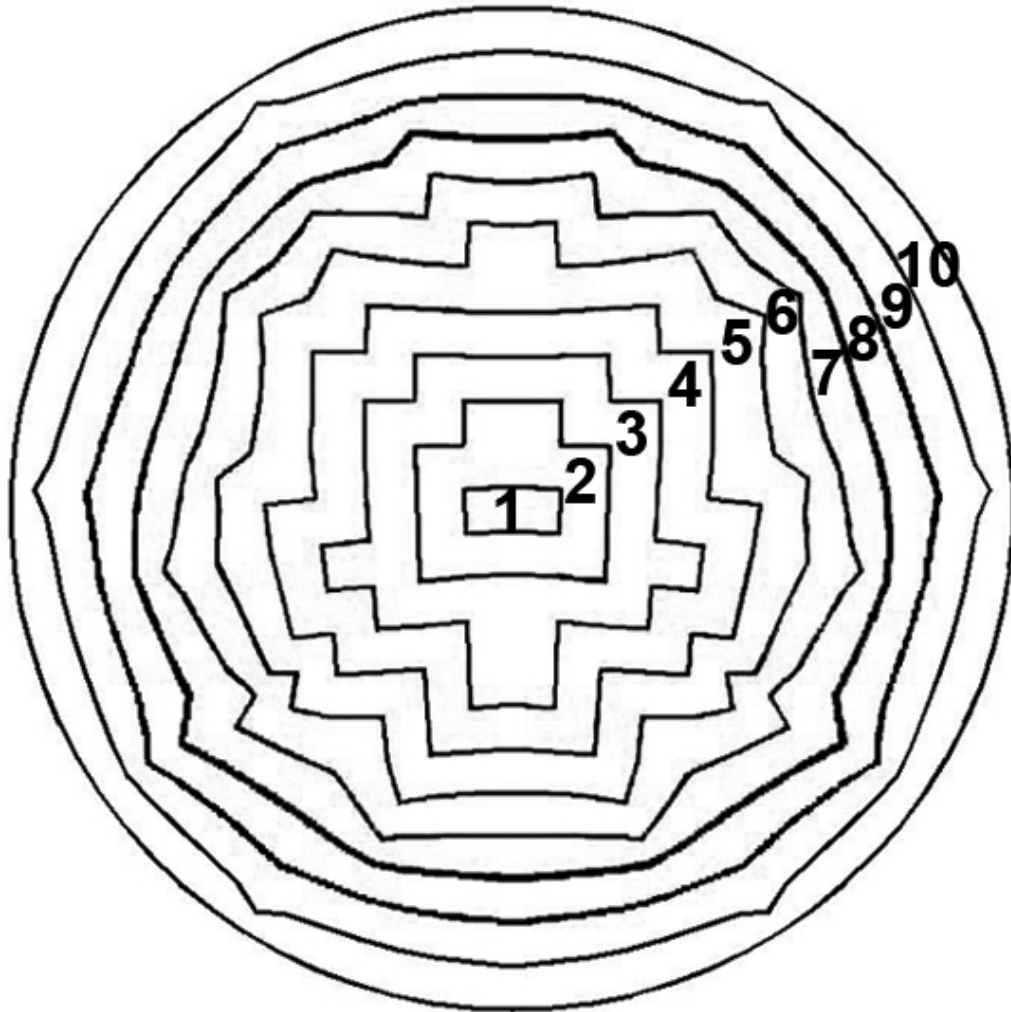
It is assumed that all the fine scales in the cell are perfectly  
 stirred reactors with a residence time  $\tau^*$

## Mechanism for H<sub>2</sub> combustion (with NO<sub>x</sub>), 13 species and 37 reactions

Units are cm<sup>3</sup>, mol, s, KJ and K

H+O <sub>2</sub> =OH+O	2.E+14	0.00	70.30	O+O+M=O <sub>2</sub> +M	6.17E+15	-0.50	0.00
H <sub>2</sub> +O=OH+H	1.8E+10	1.00	36.93	H <sub>2</sub> O <sub>2</sub> +H=H <sub>2</sub> O+OH	1.E+13	0.00	15.02
H <sub>2</sub> O+O=OH+OH	5.9E+09	1.30	71.25	H <sub>2</sub> O <sub>2</sub> +H=HO <sub>2</sub> +H <sub>2</sub>	4.79E+13	0.00	33.26
H <sub>2</sub> +OH=H <sub>2</sub> O+H	1.17E+09	1.30	15.17	O+OH+M=HO <sub>2</sub> +M	1.E+16	0.00	0.00
H+O <sub>2</sub> +M=HO <sub>2</sub> +M	2.3E+18	-0.8	0.00	H <sub>2</sub> +O <sub>2</sub> =OH+OH	1.7E+13	0.00	200.0
H <sub>2</sub> /1./ H <sub>2</sub> O/6.5/ O <sub>2</sub> /0.4/ N <sub>2</sub> /0.4/				O+N <sub>2</sub> =N+NO	1.82E+14	0.00	319.02
H+HO <sub>2</sub> =OH+OH	1.5E+14	0.00	4.20	O+NO=N+O <sub>2</sub>	3.8E+09	1.00	173.11
H+HO <sub>2</sub> =H <sub>2</sub> +O <sub>2</sub>	2.5E+13	0.00	2.93	H+NO=N+OH	2.63E+14	0.00	210.94
OH+HO <sub>2</sub> =H <sub>2</sub> O+O <sub>2</sub>	2.E+13	0.00	4.18	NO+M=N+O+M	3.98E+20	-1.50	627.65
H+H+M=H <sub>2</sub> +M	1.8E+18	-1.00	0.00	N <sub>2</sub> +M=N+N+M	3.72E+21	-1.60	941.19
H <sub>2</sub> /1./ H <sub>2</sub> O/6.5/ O <sub>2</sub> /0.4/ N <sub>2</sub> /0.4/				N <sub>2</sub> O+O=NO+NO	6.92E+13	0.00	111.41
H+OH+M=H <sub>2</sub> O+M	2.2E+22	-2.00	0.00	N <sub>2</sub> O+O=N <sub>2</sub> +O <sub>2</sub>	1.E+14	0.00	117.23
H <sub>2</sub> /1./ H <sub>2</sub> O/6.5/ O <sub>2</sub> /0.4/ N <sub>2</sub> /0.4/				N <sub>2</sub> O+N=N <sub>2</sub> +NO	1.E+13	0.00	83.14
HO <sub>2</sub> +HO <sub>2</sub> =H <sub>2</sub> O <sub>2</sub> +O <sub>2</sub>	2.E+12	0.00	0.00	N+HO <sub>2</sub> =NO+OH	1.E+13	0.00	8.31
H <sub>2</sub> O <sub>2</sub> +M=OH+OH+M	1.3E+17	0.00	190.38	N <sub>2</sub> O+H=N <sub>2</sub> +OH	7.6E+13	0.00	63.19
H <sub>2</sub> O <sub>2</sub> +OH=H <sub>2</sub> O+HO <sub>2</sub>	1.E+13	0.00	7.53	HNO+O=NO+OH	5.01E+11	0.50	8.31
O+HO <sub>2</sub> =OH+O <sub>2</sub>	2.E+13	0.00	0.00	HNO+OH=NO+H <sub>2</sub> O	1.26E+12	0.50	8.31
H+HO <sub>2</sub> =O+H <sub>2</sub> O	5.E+12	0.00	5.90	NO+HO <sub>2</sub> =HNO+O <sub>2</sub>	2.E+11	0.00	8.31
H+O+M=OH+M	6.2E+16	-0.60	0.00	HNO+HO <sub>2</sub> =NO+H <sub>2</sub> O <sub>2</sub>	3.16E+11	0.50	8.31
				HNO+H=NO+H <sub>2</sub>	1.26E+13	0.00	16.63
				HNO+M=H+NO+M	1.78E+16	0.00	203.7

GUTHEIL, E., BALAKRISHNAN, G. & WILLIAMS, F. A. (1993) Structure and extinction of hydrogen-air diffusion flames. IN PETERS, N. & ROGG, B. (Eds.) *Reduced kinetic mechanisms for applications in combustion systems*. New York, Springer-Verlag



Burst disk opening process was approximated by a 10 step process.

Sections were open from 1 to 10 at times calculated according to the table in the next slide.

# Burst disk opening time

Following Spence & Woods (1964), the following formula was used to obtain the rupture time of a burst disk:

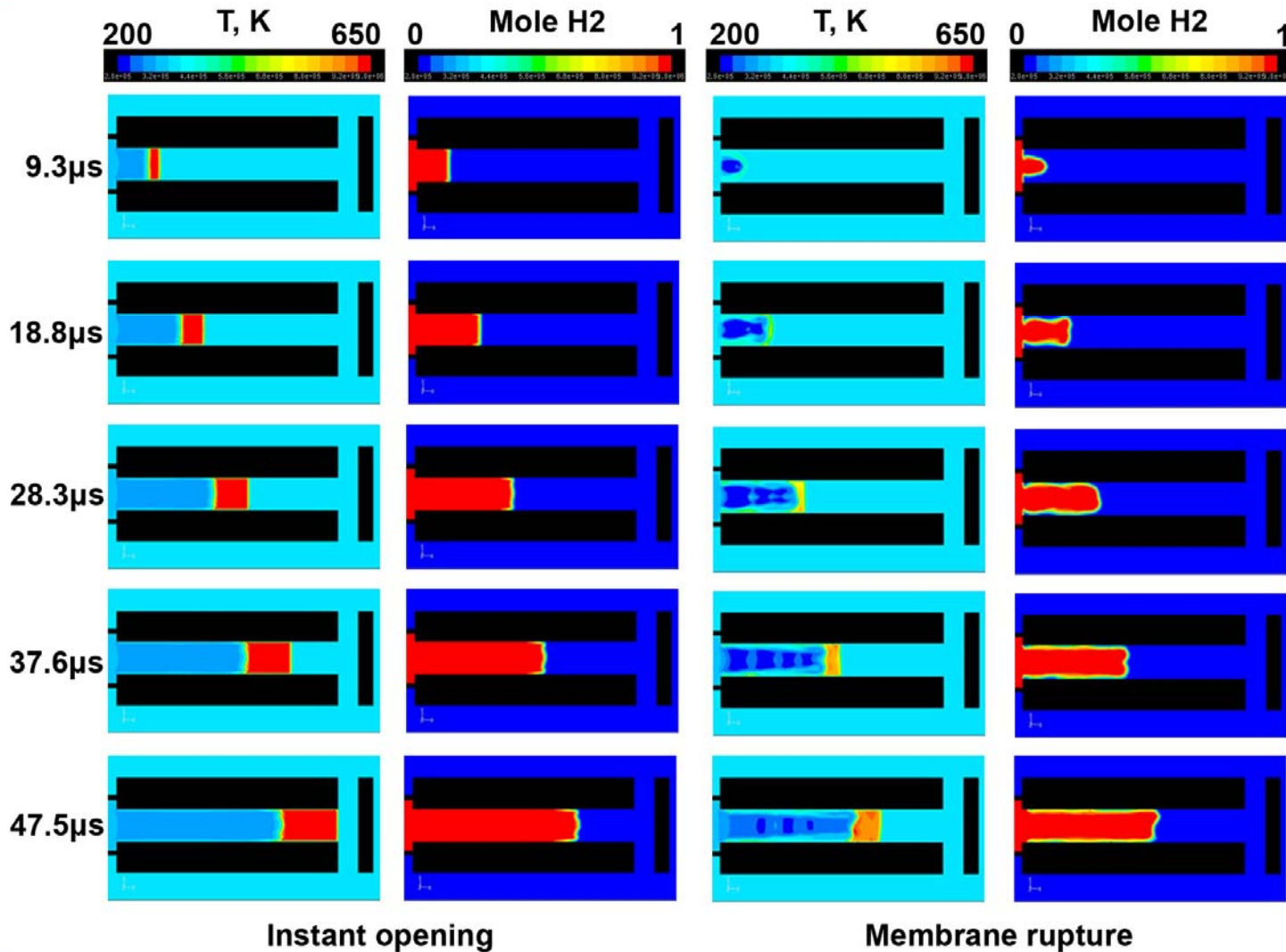
$$t = k \left( \frac{\rho b d}{p} \right)^{1/2}$$

where  $\rho$  – the density of the diaphragm material (annealed copper),  
 $b$  and  $d$  are the thickness and the diameter of the diaphragm,  
 $k$  – constant, 0.92  
 $p$  – driving pressure

Opening times of sections for each of the simulations are listed below:

Section	1	2	3	4	5	6	7	8	9	10
Opening time, $\mu\text{s}$ [1.35 MPa case]	0	4.7	9.4	14.2	18.9	23.6	28.4	33.1	37.8	42.6
Opening time, $\mu\text{s}$ [1.65 MPa case]	0	4.3	8.6	12.8	17.1	21.4	25.6	29.9	34.3	38.5
Opening time, $\mu\text{s}$ [2.43 MPa case]	0	3.5	7.1	10.6	14.2	17.7	21.3	24.8	28.3	31.9
Opening time, $\mu\text{s}$ [2.9 MPa case]	0	3.2	6.5	9.7	12.9	16.2	19.4	22.6	25.9	29.1

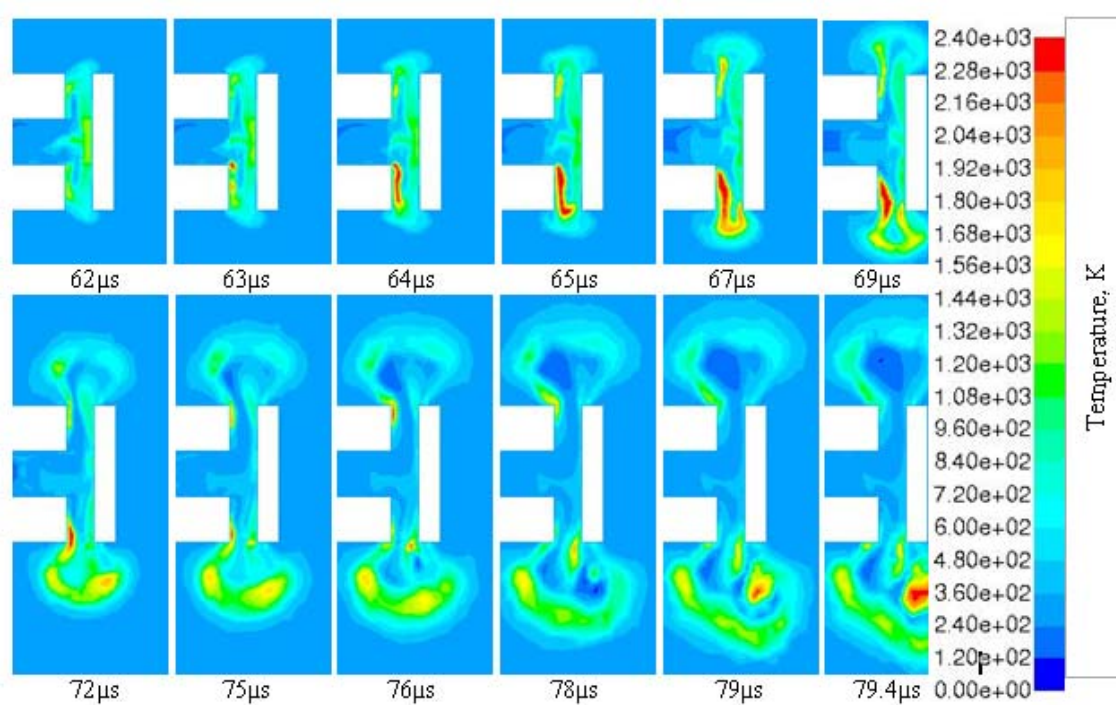
# Instant vs finite opening time



# Simulation results



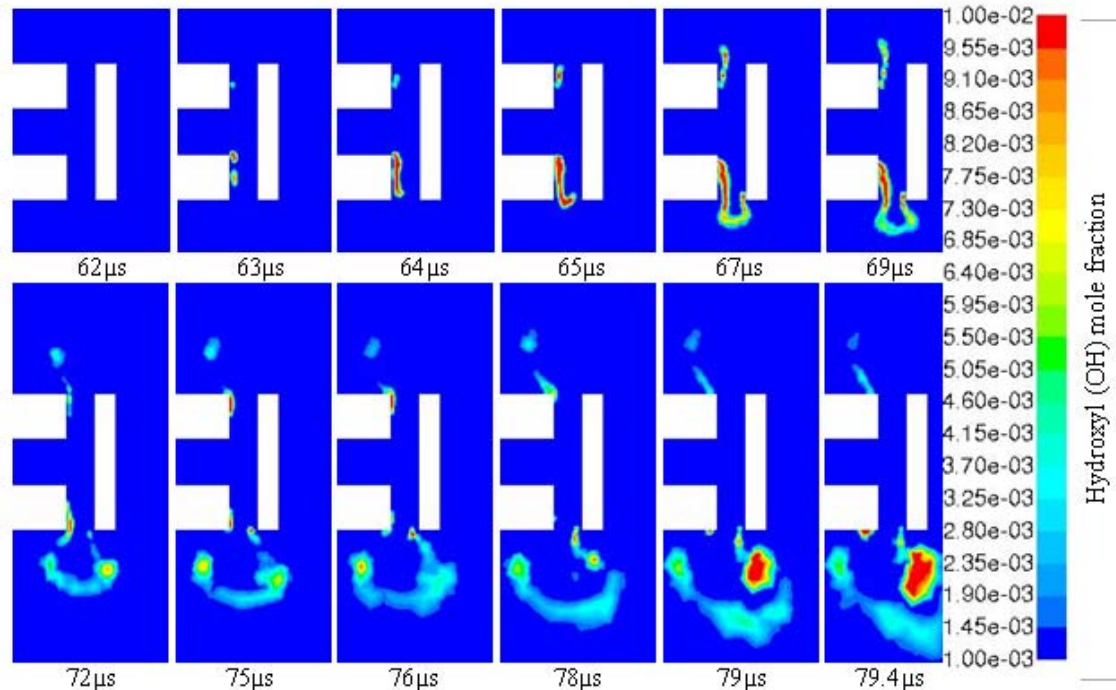
# P=29 bar



Ignition initiated in the radial vent channel and extinguished.

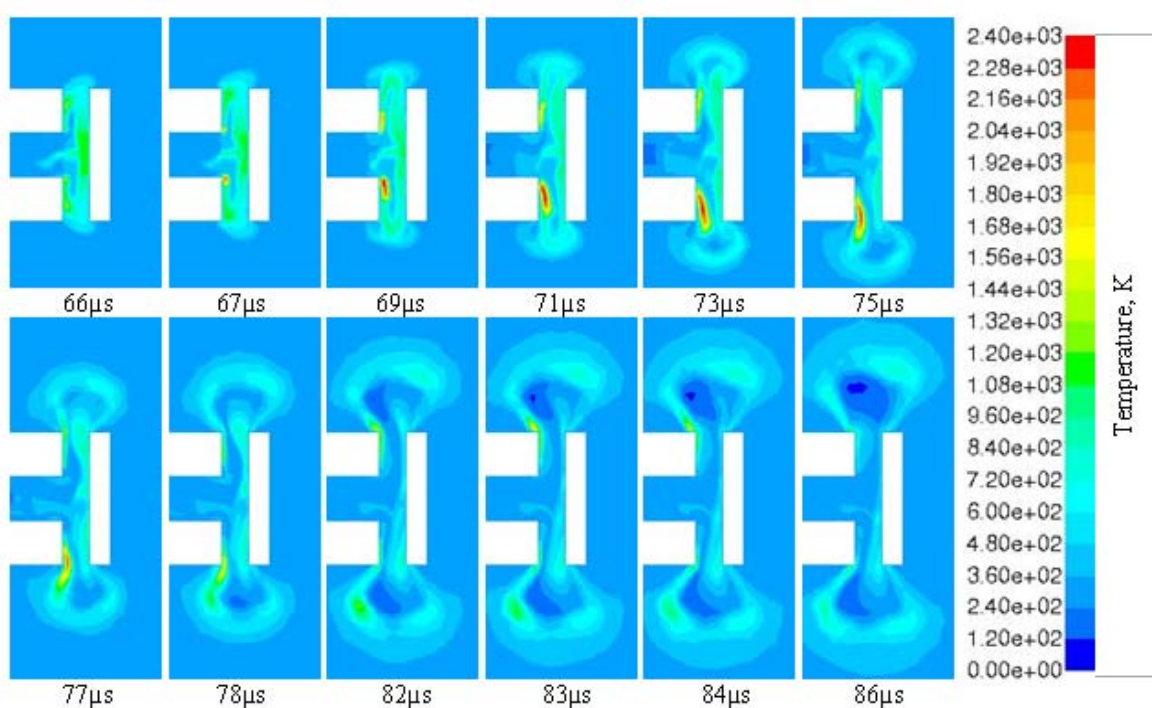
Combustion is reinitiated in a number of spots outside the PRD.

Concentration of hydrogen in these spots just before the ignition is in the range 29-36% by vol.





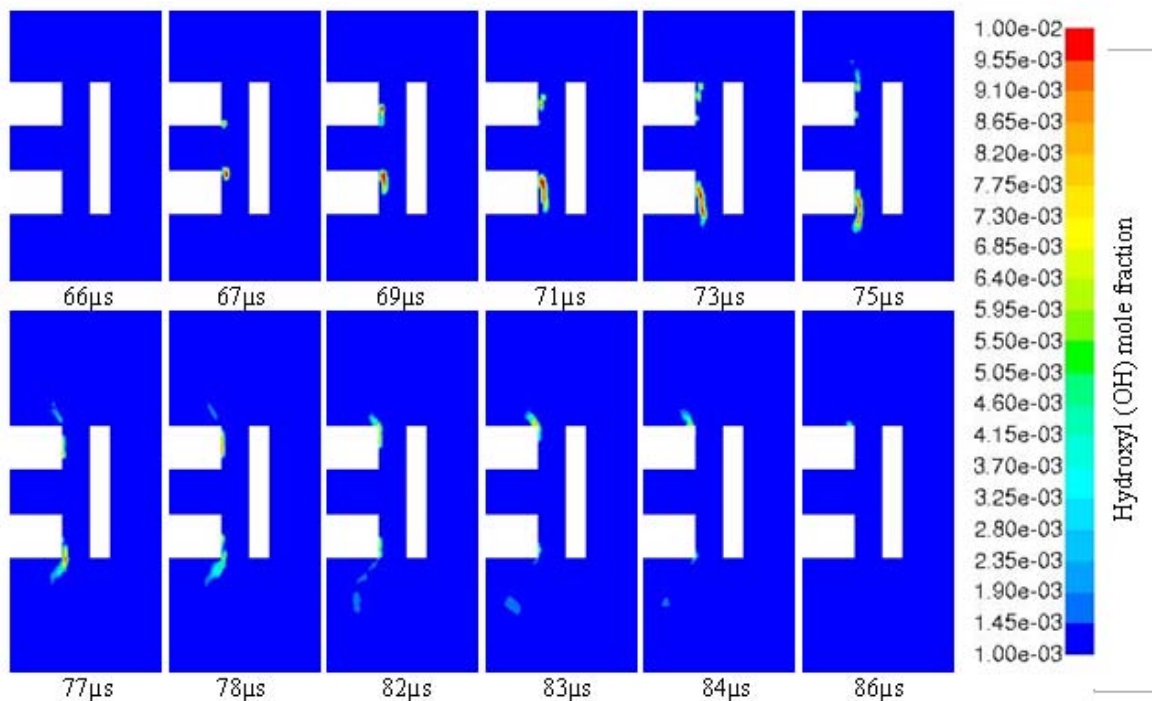
# P=24.3 bar



Ignition is initiated in the radial channel (closer to the upstream wall – second shock reflection) similar to 29 bar case

Reignition outside is not observed.

Process took 10 $\mu$ s and was experimentally registered (“ignition” in private communication)



**P=16.5 bar**

## Temperature

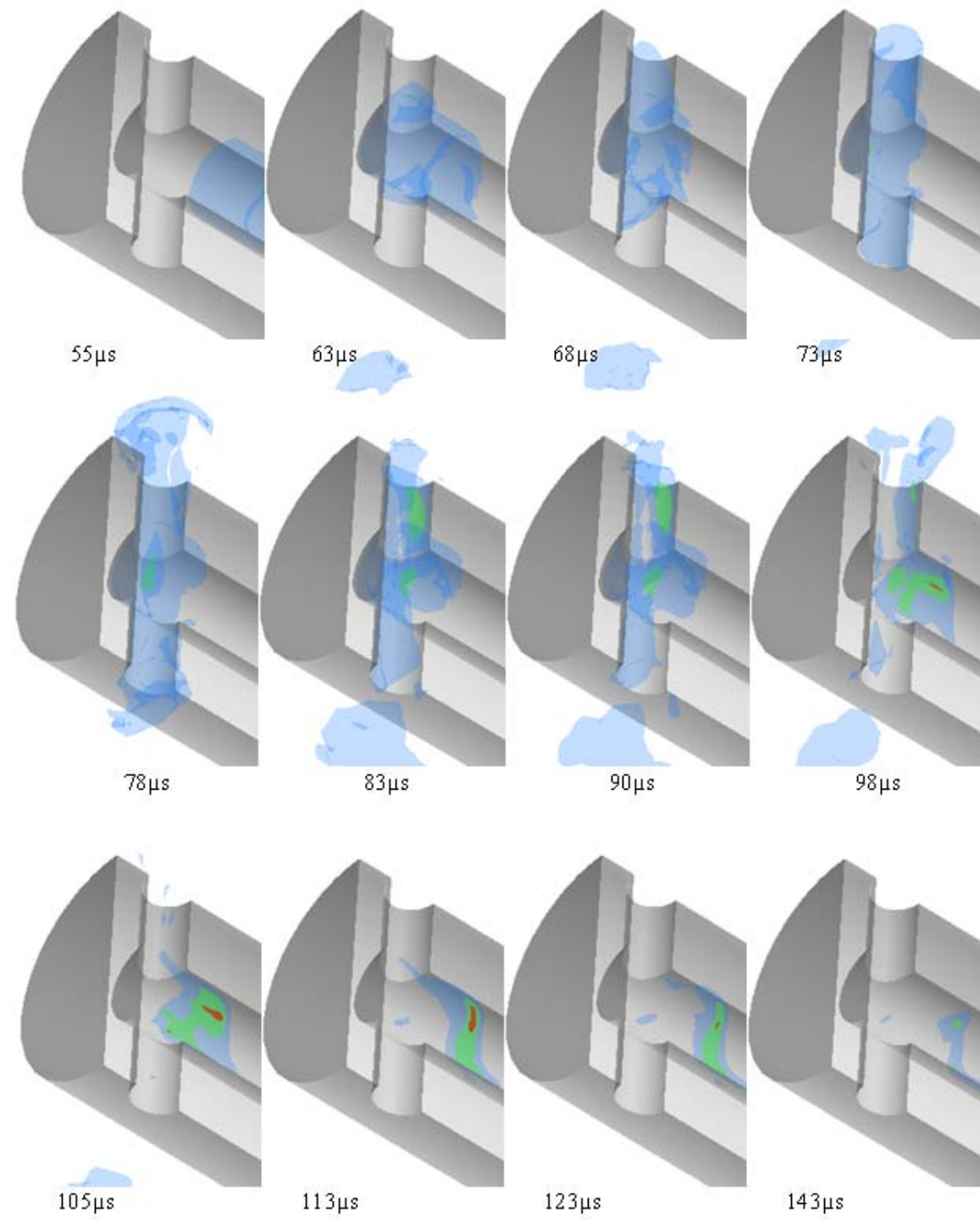
Ignition spot is located in the axial channel (not in a symmetry plane). The combustion propagates upstream (until oxygen is consumed)!

Iso-surfaces are:

Blue – 550 K,

Green – 1500 K,

Red – 2400 K



**P=16.5 bar**

## **OH mole fraction**

Ignition is confirmed by a hydroxyl mole fraction distribution.

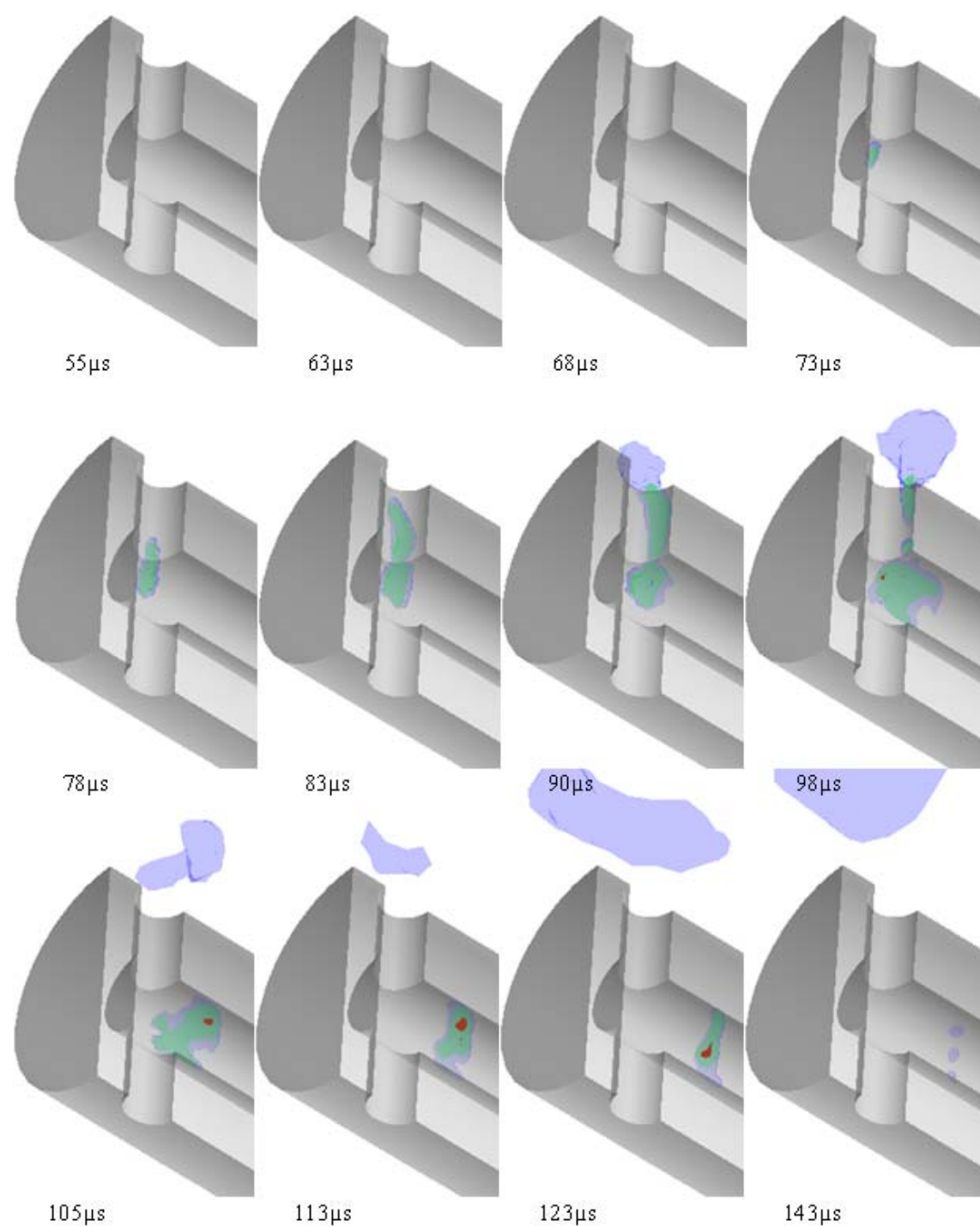
Light sensor in experiments was located outside and could not have registered ignition.

Iso-surfaces are:

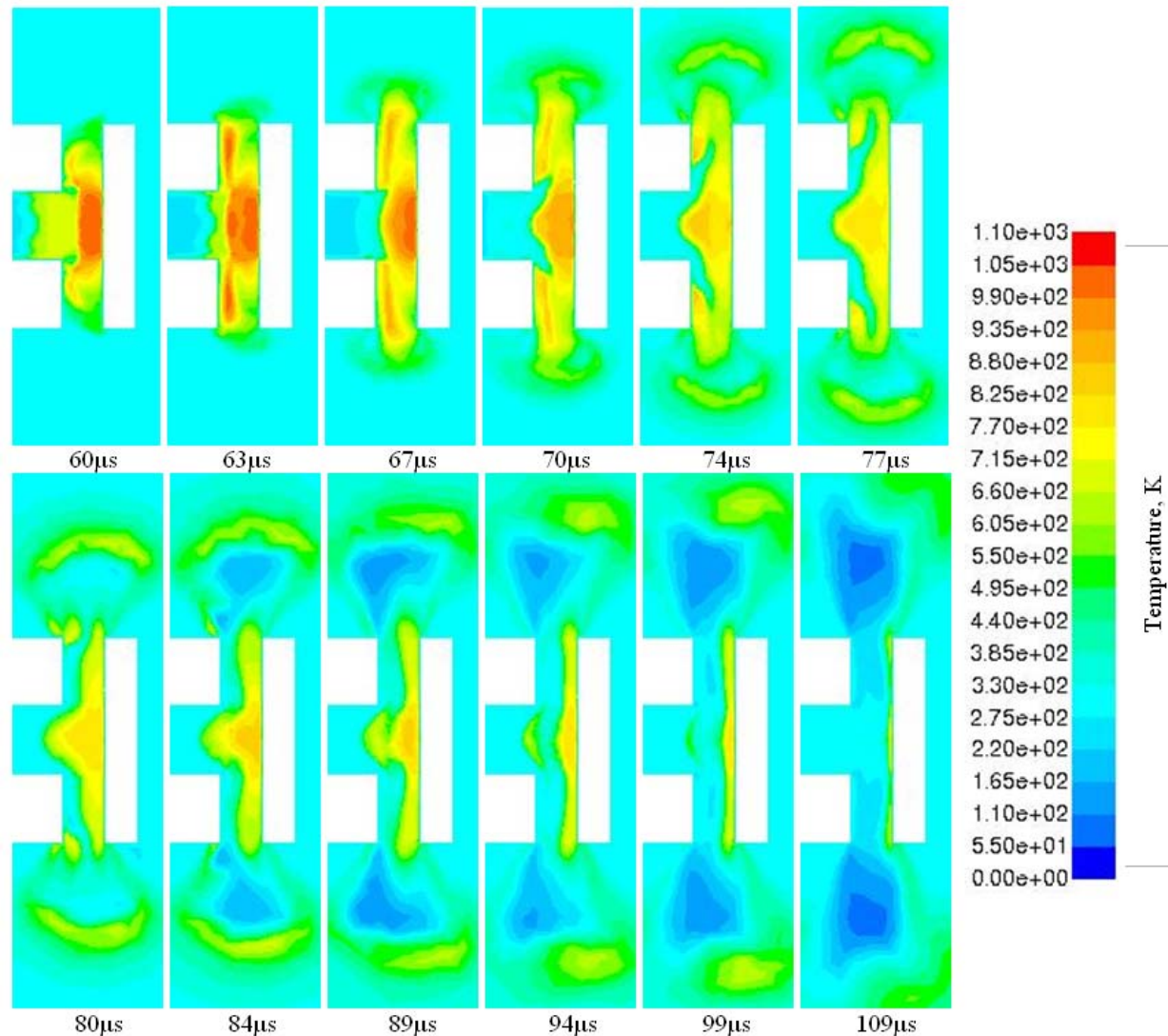
Blue – 0.0002,

Green – 0.002,

Red – 0.02



# P=13.5 bar – no ignition



Hydroxyl mole fraction  
 is  $\sim 10^{-9}$ - $10^{-10}$   
 (not shown).

Note: temperature  
 scale is different in  
 this picture.



# Conclusions

- The LES model based on EDC with detailed kinetics and inertial burst disk was successfully applied to reproduce and explain experimental observations in T-shape channel (PRD mock-up).
- The lower pressure limit for spontaneous ignition in a T-shaped PRD is established: there is no ignition at storage pressure 13.5 bar; there is ignition followed by self-extinction at pressures 16.5 and 24.3 bar; and there is “sustainable” ignition (reignition outside the channel) for storage pressure 29 bar.
- The ignition is located in the area of the second reflection of initial shock and where the concentration of hydrogen in air is in the range 29-36% vol.
- The model can be used as a contemporary tool for design of innovative pressure relief devices.



MSc in Hydrogen Safety Engineering (distance learning course):  
<http://campusone.ulster.ac.uk/potential/postgraduate.php?ppid=24>